# Study of UV fluorescences induced from $4f^3 \rightarrow 4f^25d$ multistep absorptions of Nd<sup>3+</sup>ions in YLiF<sub>4</sub> and LuLiF<sub>4</sub> crystals <sup>1</sup>

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Nd<sup>3+</sup> ultraviolet (UV) fluorescence induced by multiphotonic laser excitations was studied in doped Nd:YLiF, (YLF) and Nd:LuLiF<sub>4</sub>(LLF) crystals by using the time resolved spectroscopy technique. The UV luminescences are due to transitions between the 4f<sup>2</sup>5d and the 4f<sup>3</sup> electronic configurations of Nd<sup>3+</sup> ions. The 4f<sup>2</sup>5d configuration can be reached by direct pumping or by multiphotonic excitation, both processes give rise to the UV band emission with structure due to the strong phonon coupling expected for 5d orbital involvement in the transition. The multiphotonic excitation process is due to three photons (532 nanometers [nm]) sequential absorptions by metastable levels of the 4f<sup>3</sup> configuration split by crystalline local field. The sequential excitation of Nd by the laser excitation is attributed to the  $4I_{9/2}$  + 532 nm  $\rightarrow 4G_{7/2}$  ground state absorption followed by the  $4G_{7/2}$  + 532 nm  $\rightarrow 2F_{5/2}$  and  $2F_{5/2}$ + 532 nm  $\rightarrow$  4f<sup>2</sup>5d excited state absorptions. The UV emissions due to 4f<sup>2</sup>5d configuration are parity allowed, having lifetime of 35 nanoseconds (ns) in contrast to UV emissions from 4f<sup>3</sup> configuration which are induced by two absorption steps and are parity forbidden showing longer lifetime of 8 microseconds  $(\mu s)$  and narrow lines. The polarization effects of the UV emissions were studied and their behaviors are dependent on the excited state configuration involving or not involving the 5d orbital. The allowed UV emission positions were affected by the host variation more than the ones originated from the  $4f^3$  configuration as expected. The electronic energy of the 4f<sup>2</sup>5d configuration shifts to lower energy for increasing the crystal field.

> Key words: Multiphotonic excitation process. Ultraviolet radiation. UV transition.

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## 1 Introduction

Rare earth ions-doped crystals are very useful laser media for generating laser radiation in the visible and infrared region. Solid state materials doped with Nd are very promising for using as laser medium for ultraviolet (UV) region (THO-GERSEN; GILL; HAUGEN, 1996). The excitation can be performed directly to the level of interest or sequentially pumped. The most interesting pumping mechanism is the three steps excitations at 532 nanometers (nm) which has the advantage of matching the second harmonic of Nd:YAG laser which is one of the most disseminated laser for optical pumping systems (THOGERSEN; GILL; HAUGEN, 1996; VENIKOUAS et al., 1984). The  $4f^n \rightarrow 4f^{n-1}5d$  transitions are characterized by a strong environmental interaction and they are responsible for high oscillator strength and broad band absorption and emission spectra in the UV range. Otherwise, the intraconfigurational  $4f^n \rightarrow$ 4f<sup>n</sup> transitions are parity forbidden and they are sharp and weak because they take place only due to the crystal field configuration mixing (KOLLIA et al., 1998; POWELL, 1998<sup>1</sup>).

The spectrum characterization of the  $4f^3 \rightarrow 4f^25d$  and  $4f^3 \rightarrow 4f^3$  transitions provides information about the local level structure and electronphonon coupling differences between  $4f^3$  and the  $4f^25d$  configurations (KOECHNER, 1986).

The pumping scheme used in this work is illustrated in Graphic 1. The main UV emissions originated form the 4f<sup>2</sup>5d configuration are indicated in this illustration.

# 2 Experimental setup

YLF and LLF samples were cut and polished properly with the c-axis parallel to the longest side of the rectangular samples. The absorption



Source: Toghersen and collaborators (1996).

spectra were performed using a Cary-Olis 17 D double-beam spectrophotometer interfaced to a computer. Using a time resolved spectroscopy system of 10 nanoseconds (ns) of resolution the emission spectra and decay time determination were provided. The laser pumping system consists of a frequency doubled Nd:YAG pulsed laser whose beam intensity is reduced and focused on the samples. The emission of the samples is focused into the monochromator that disperses and directs the light to the EMI S-20 photomultiplier tube. The detection system is connected to the 200 megahertz (MHz) Tektronix oscilloscope and a Box-Car gated integrator coupled to a microcomputer as showed in the Illustration 1.





Source: The authors.

The reduction of laser energies to hundred of microjoules ( $\mu$ J) is necessary to avoid the self focusing and thermal lens problems in the millijoule (mJ) energy range which can destroy the samples (KOLLIA et al., 1998). The thermal lens effects produce a strong signal decrease and a bad signal-noise ratio, which can disguise the UV emission signal. A non-divergent beam and energies ranging from 4 to 100  $\mu$ J can be used to improve UV fluorescence performance.

#### 3 Results and discussion

In order to obtain a time resolved spectroscopy of the emissions and to investigate polarization effects they were measured both UV emissions from 4f<sup>3</sup> and 4f<sup>2</sup>5d excited states, which follow in the sub-sections.

### 3.1 UV Emissions and the Polarization Effects

The frequency doubled Nd:YAG-Qswitched laser operating at 532 nm with a repetition rate of 10 hertz (Hz) was used in all the investigations of UV fluorescences. The lower energy limit of the 4f<sup>2</sup>5d band of Nd<sup>3+</sup> in YLF lies at ~ 55,000 cm<sup>-1</sup> so, the three-photon sequential absorptions of 532 nm, can excite Nd ion to this configuration. The UV fluorescence spectra of Nd in YLF crystal are exhibited in Illustrations 3 and 4. The emission spectrum of  $4f^25d \rightarrow 2G_{_{7/2}}$  at 230 nm and  $4f^25d \rightarrow 4F_{_{9/2}}$  at 260 nm are exhibited by Illustration 3. They were discriminated by using a narrow static gate (2 ns) and a delay time of 20 ns, in order to detect only the allowed UV emissions. The two-photon that induced emissions from 4f<sup>3</sup> configuration are also present in the experiment and were optically discriminated by using a narrow static gate of 2 ns using a longer delay time of 2 microseconds ( $\mu$ s) in the box-car averager. By the comparison of both integrated emission signals from distinct electronic configurations, we must say that UV emission from 4f<sup>2</sup>5d is about ten times stronger than the one from 4f<sup>3</sup> configuration, besides having an absorption process of order two. The emission spectrum of Illustration 4 exhibits the following transitions:  $2F_{5/2} \rightarrow 4I_{15/2}$ at 310 nm,  $2F_{5/2} \rightarrow 4I_{13/2}$  at 290 nm,  $2F_{5/2} \rightarrow 4I_{11/2}$ at 275 nm and  $2F_{5/2} \rightarrow 4I_{9/2}$  at 260 nm. The measured lifetime of  $2F_{5/2}$  state is 8  $\mu$ s in contrast with the measured lifetime of 35 ns found for the 4f<sup>2</sup>5d configuration. By comparing these results one concludes that the 4f<sup>2</sup>5d configuration has an emission spectra larger than the emission from  $4f^{3}(2F_{5/2})$  state. This comes from the fact that 4f<sup>3</sup> configuration produces inner states shielded by the closed 5p<sup>6</sup> orbital thus producing optical transitions with very small phonon coupling. As a consequence, the polarization effects in the UV emissions from each configuration must have different effects. The strongest polarization effects are observed for the internal 4f<sup>3</sup> transitions, while the 4f<sup>2</sup>5d-mixed configuration is slightly polarized.



Graphic 2: Polarized emission spectrum of the 4f<sup>2</sup>5d configuration of Nd:YLF crystals, after three photons absorptions of 532 nm





Source: The authors.

#### 3.2 Host lattice effects

UV emissions were induced in both Nddoped crystals of YLF and LLF with laser pulsed excitation at 532 nm. Graphic 4 shows the measured 4f<sup>2</sup>5d-emission spectrum of both crystals for comparison. The spectrum shape of both UV emissions is similar. However, a strong shift of main peaks by 180 and 200 cm<sup>-1</sup>, towards the lowest energy, was observed for LLF crystal. This suggests that the total energy of the bottom of the 4f<sup>2</sup>5d mixed configuration of Nd ions in LLF crystal is diminished by approximately 180-200 cm<sup>-1</sup> in comparison with the position of this level in YLF, due to the strongest crystal field felt by Nd ions in LLF crystal when Y<sup>3+</sup> neighbors are substituted by Lu<sup>3+</sup> ions which have smallest ionic radius. On the other hand, the 4f<sup>3</sup> configuration has UV emissions showing a very small shifting caused by the local field increasing. As observed, the  $2F_{5/2} \rightarrow 4I_{9/2}$  emission showed a 20 cm<sup>-1</sup> peak shift.





Source: The authors.



mechanism of the  $4f^3$  configuration. The fit is S = cE<sup>1.8</sup>

Source: The authors.



Graphics 5 and 6 show the behavior of the 260 nm UV emissions intensities *versus* the laser pumping energy at 532 nm. Illustration 6 a shows the behavior of  $2F_{5/2}$  state emission at 260 nm from the 4f<sup>3</sup> configuration, giving an energy power dependence of second order (c = 1,8) for a pumping energy varying from 30 to 200  $\mu$ J. The UV emission at 260 nm from the 4f<sup>2</sup>5d configuration exhibits power law dependence according to the three order process (Graphic 7). In this case, it is seen a luminescence decrease for pumping energy higher than 130  $\mu$ J.

#### 4 Final considerations

The Nd:LLF crystal shows a three photon multistep absorption at 532 nm allowing the population of a  $4f^25d$  configuration, similar as it was previously observed in Nd:YLF crystal. This

observation validates the use of LLF crystal as a promising system for UV laser operation near 260 nm. The spectral discrimination of a long lived UV emission of 4f<sup>3</sup> configuration from the UV fast emission component allowed us to clarify the multistep excitation mechanism involved in neodymium UV fluorescence for laser applications.

#### **Notes**

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- 2 See chapter 10.
- 3 This transformation is related to the equation  $c = \lambda x f$ , where c is the light speed,  $\lambda$  is the wavelength and f is the frequency that is proportional to the photon energy.

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